



IEA/RFP-WS Feb. 29, 2000, U. Wisc.

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Understanding of Mode-locking Phenomena in TPE-RX

Recent measurements of the vessel current, mode rotation in the B_t signals, T_{e0} and n_{e0} give a better picture to understand the threshold for the mode to lock, in line with the Fitzpatrick's theory.

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**(Acknowledgment to R. Fitzpatrick², T. Bolzonella³,
P. Zanca³ and S. Martini³, 2: U. Texas, 3:
Consorzio RFX)**



Outline

- Introduction
- Locked (LM) and nonlocked (NonLM) discharges
- Diagnostics (B_r and B_t -arrays)
- B_t signals
 - *Raw data, mode amplitude, velocity and spectrum*
- Comparison of T_{e0} , T_i , n_{el}^* , β_p and τ_E
- Comparison with Fitzpatrick's threshold
- Summary



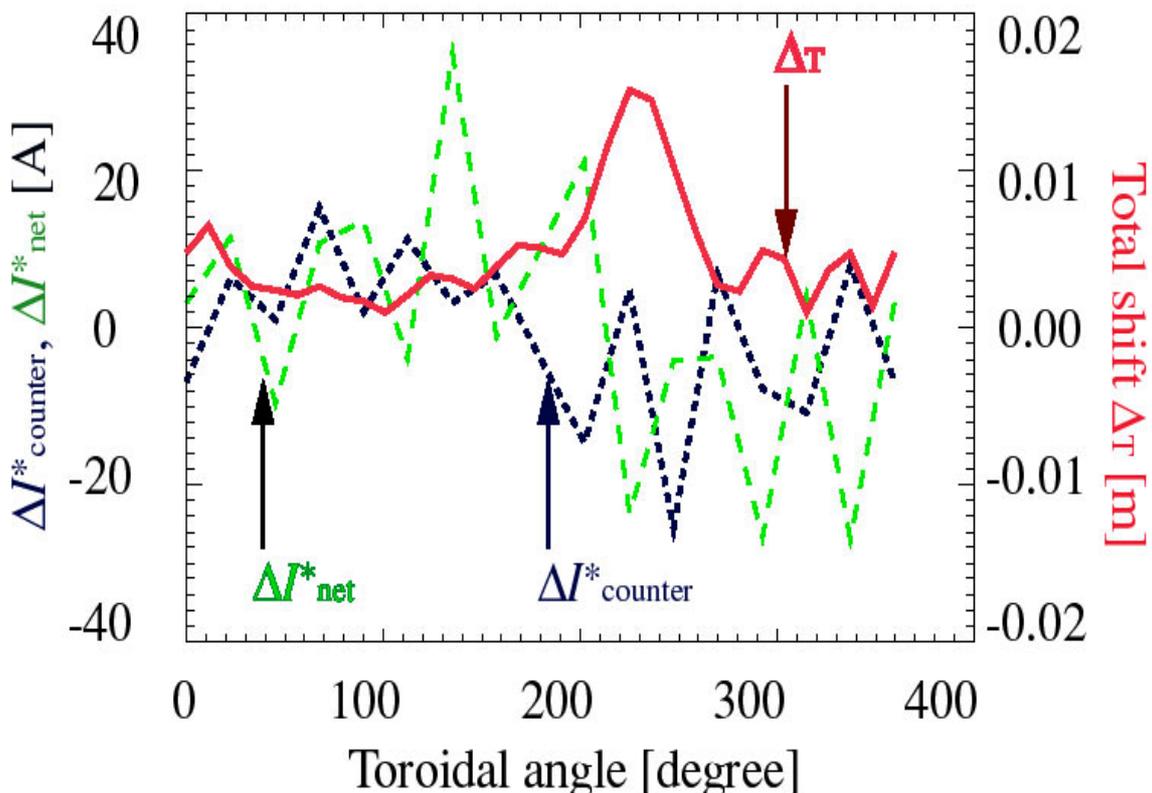
Introduction

- Phase- and wall-locked mode (**LM**) has been observed in many RFPs (RFX, T2, MST, TPE-RX).
- In TPE-RX, typical LM discharges as well as the discharges in which 'slinky' structure disappears exist (**NonLM**) [Phys. Plasmas 6 (1999) 3824].
- *Two hypotheses* for the cause of the mode locking have been experimentally checked.
 - (1) *Halo current* => O(30-40A).
Not the cause but the result of the shift of the plasma column
[Koguchi et al., in preparation].
 - (2) *Fitzpatrick's theory* [Phs. Plasmas 6 (1999) 3878] =>
Investigated here.



Vessel Current Measurement

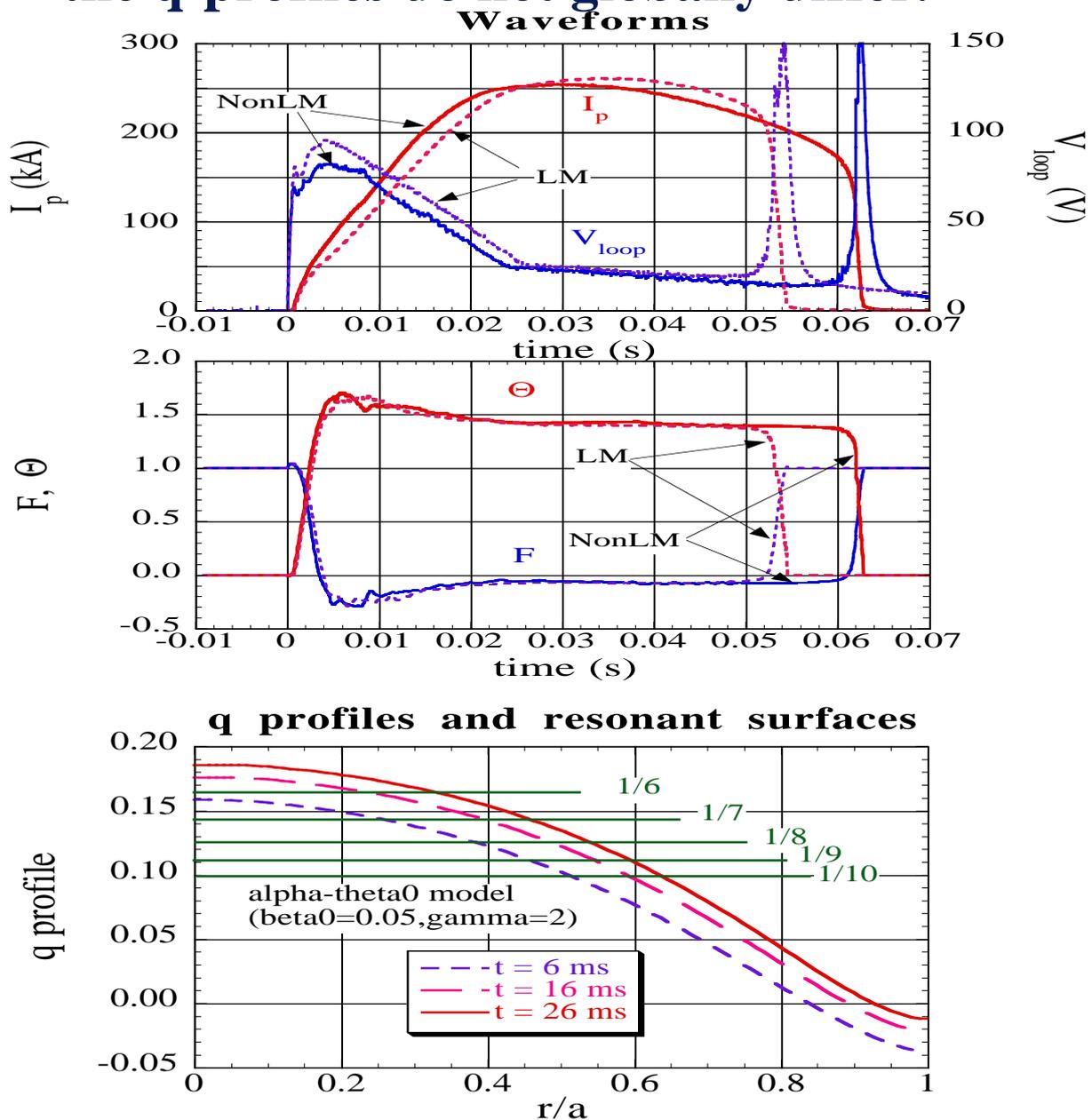
- Toroidal distribution of the vessel current was directly measured in order to check if a significant amount of the halo current is formed.
- The result showed that **the vessel current is 30-40 A**, which is much weaker than what is expected to form the halo current, and is an order of the ion saturation current as a result of the local shift of the plasma column.





LM and NonLM Discharges

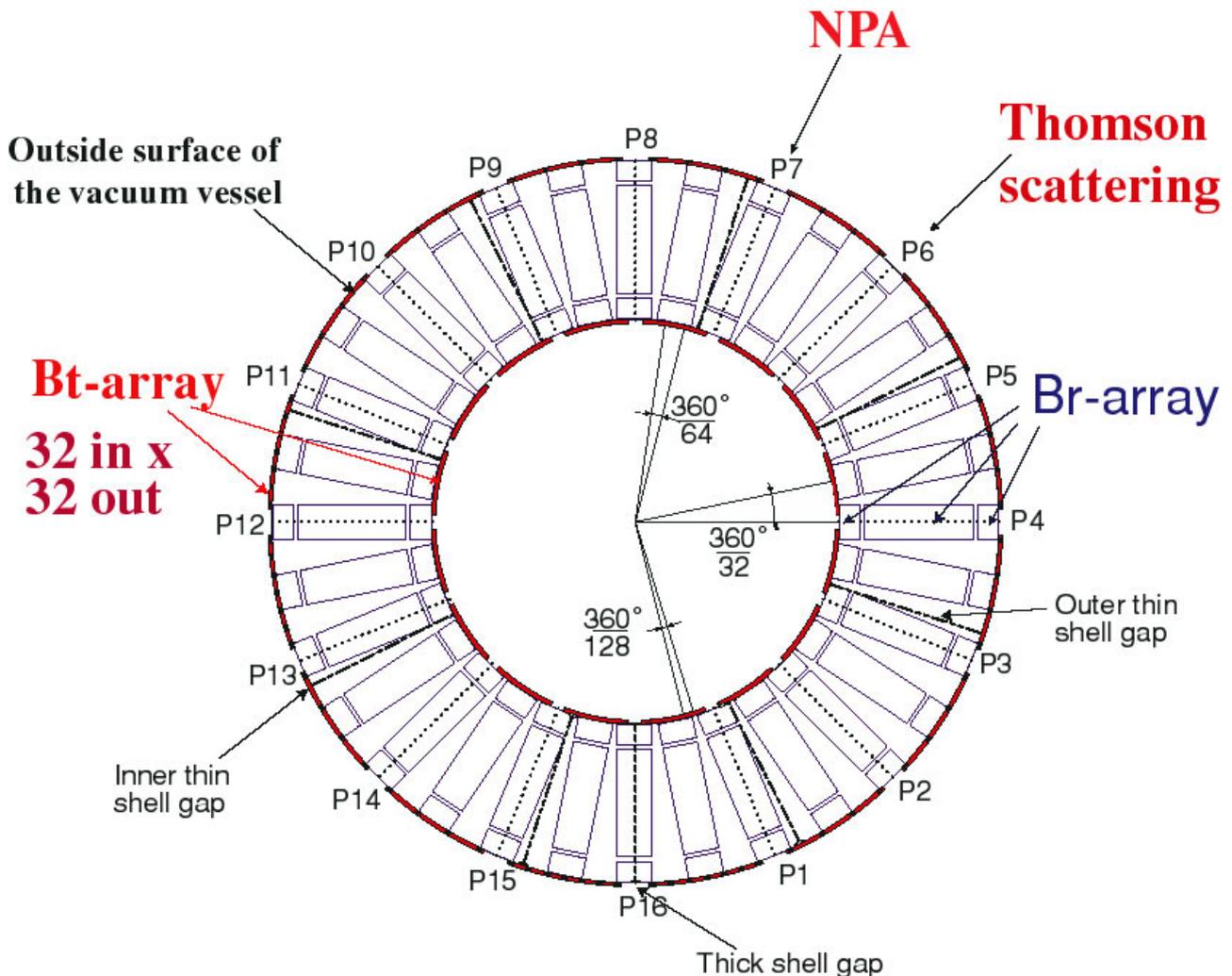
- LM and NonLM is compared in detail at $I_p = 250$ kA, $\Theta = 1.4$ and $F = -0.1$.
- Note that values of (F, Θ) are almost the same between LM and NonLM so that the q -profiles do not globally differ.





Diagnositics

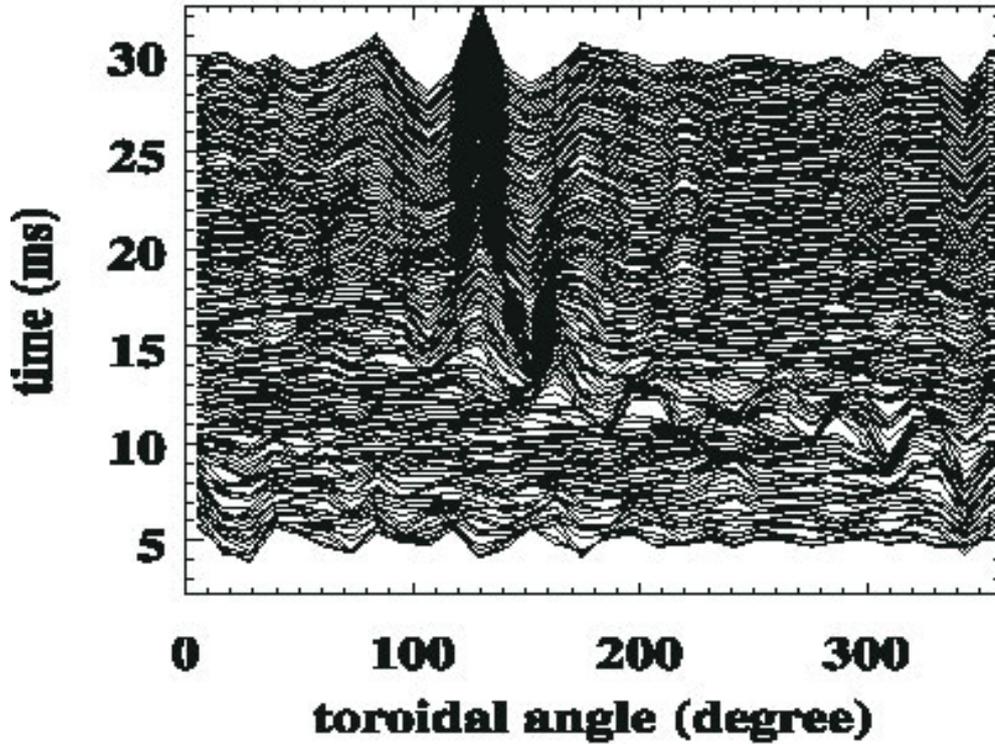
- **Br** => good for the reconstruction of the LCFS [P. Zanca and S. Martini, Plasma phys. Control. Fusion 41 (1999) 1251].
- **Bt** => larger than Br at the edge and good for studies on MHD activities in general (including mode rotation).



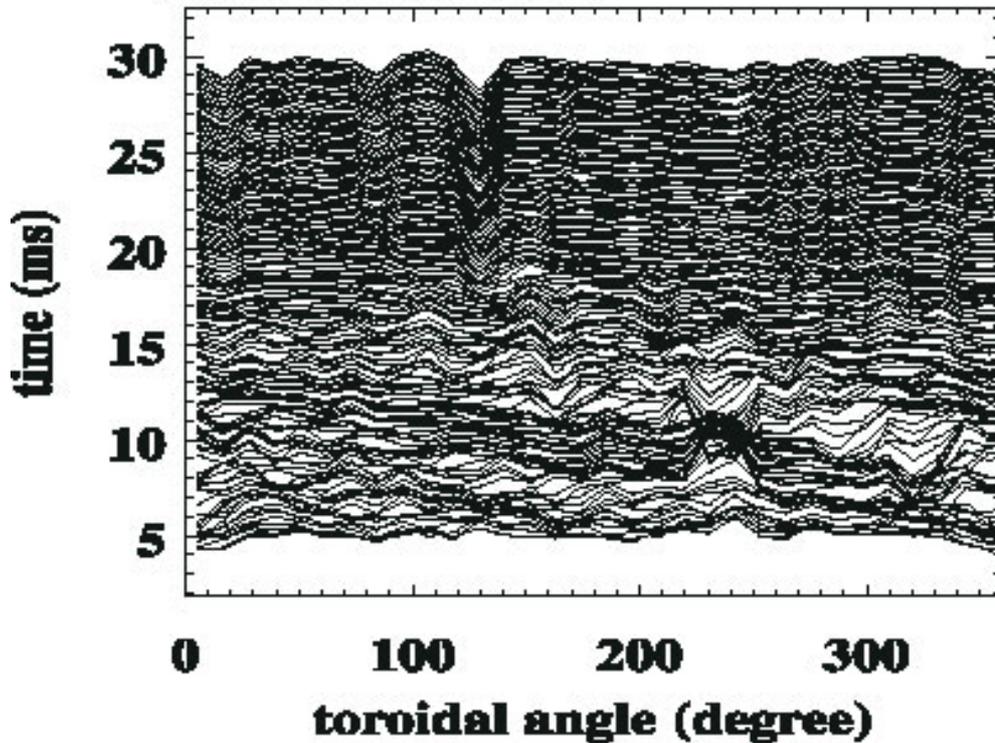


Bt Array Signals in LM

M1 : Z=0.0299 T



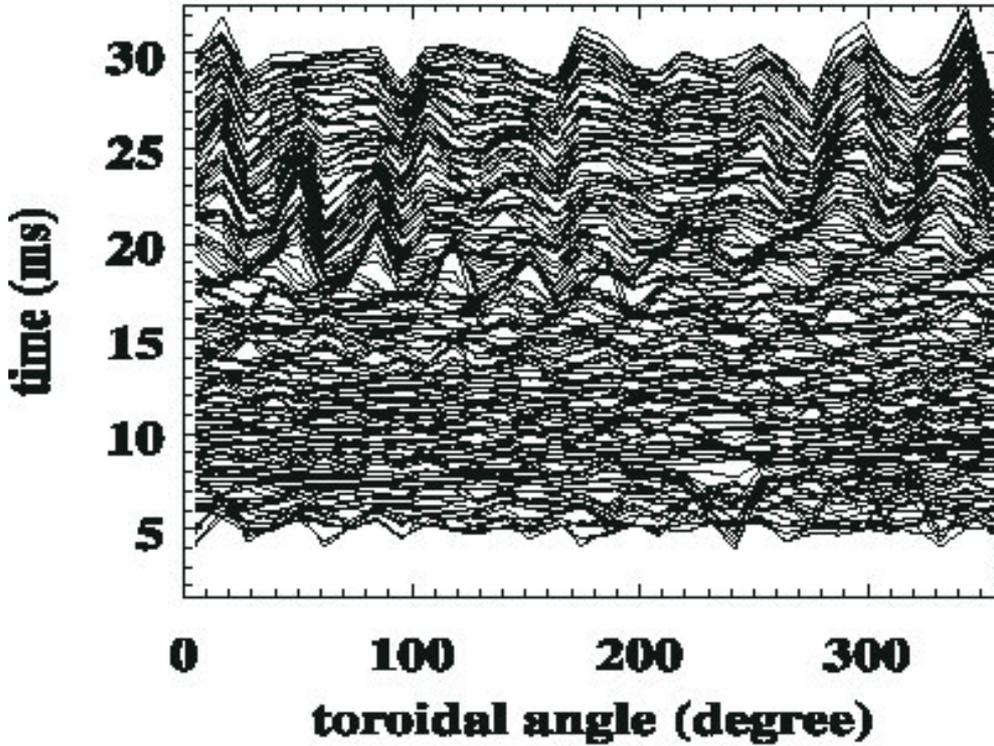
M0 : Z=0.0122 T



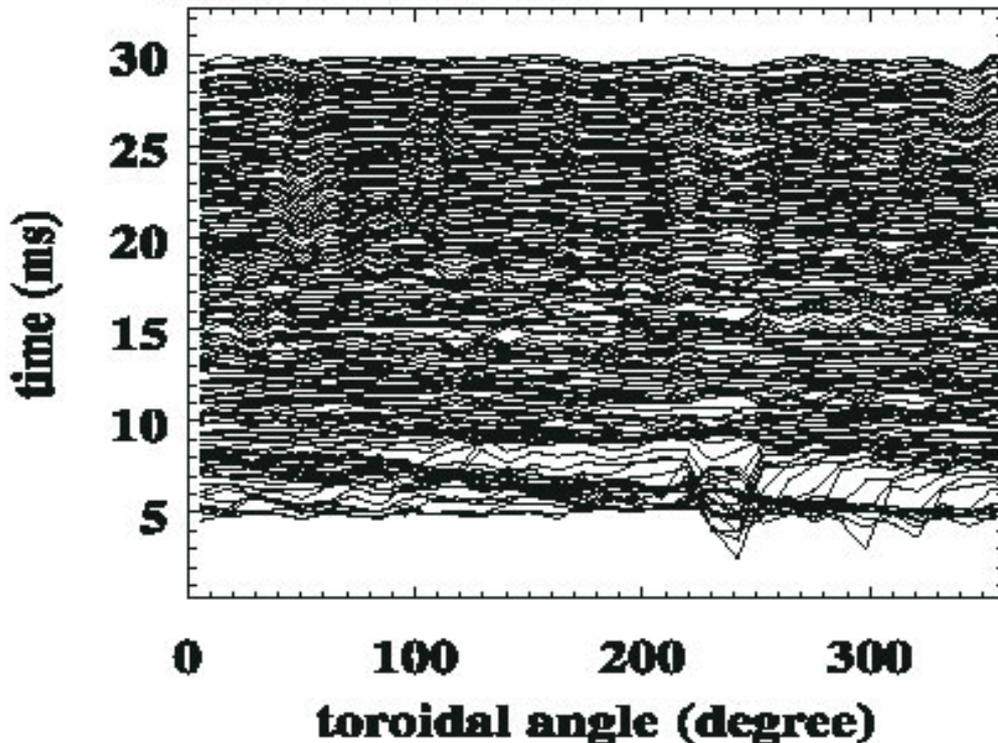


Bt Array Signals in **NonLM**

M1 : Z=0.0132 T

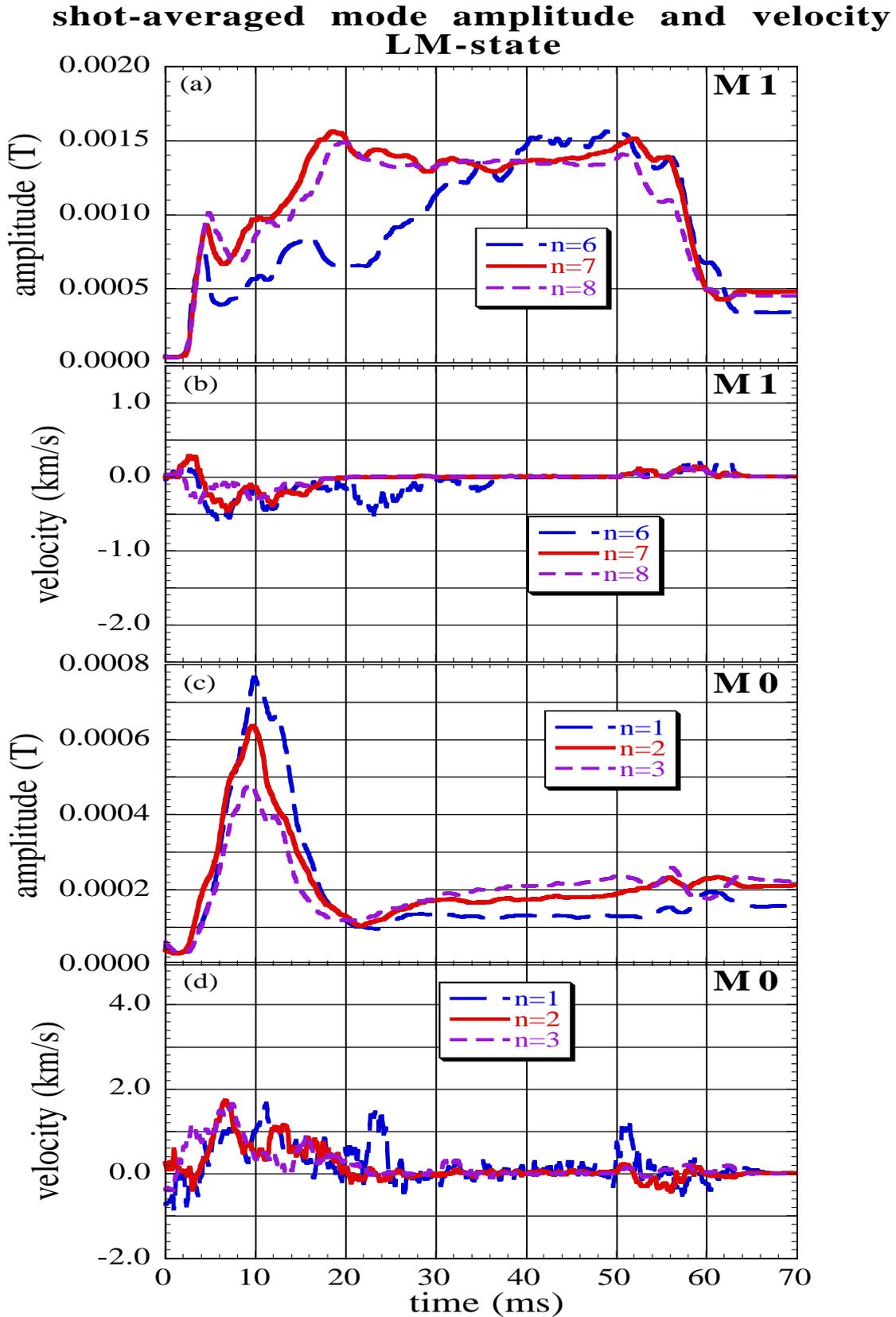


M0 : Z=0.0098 T



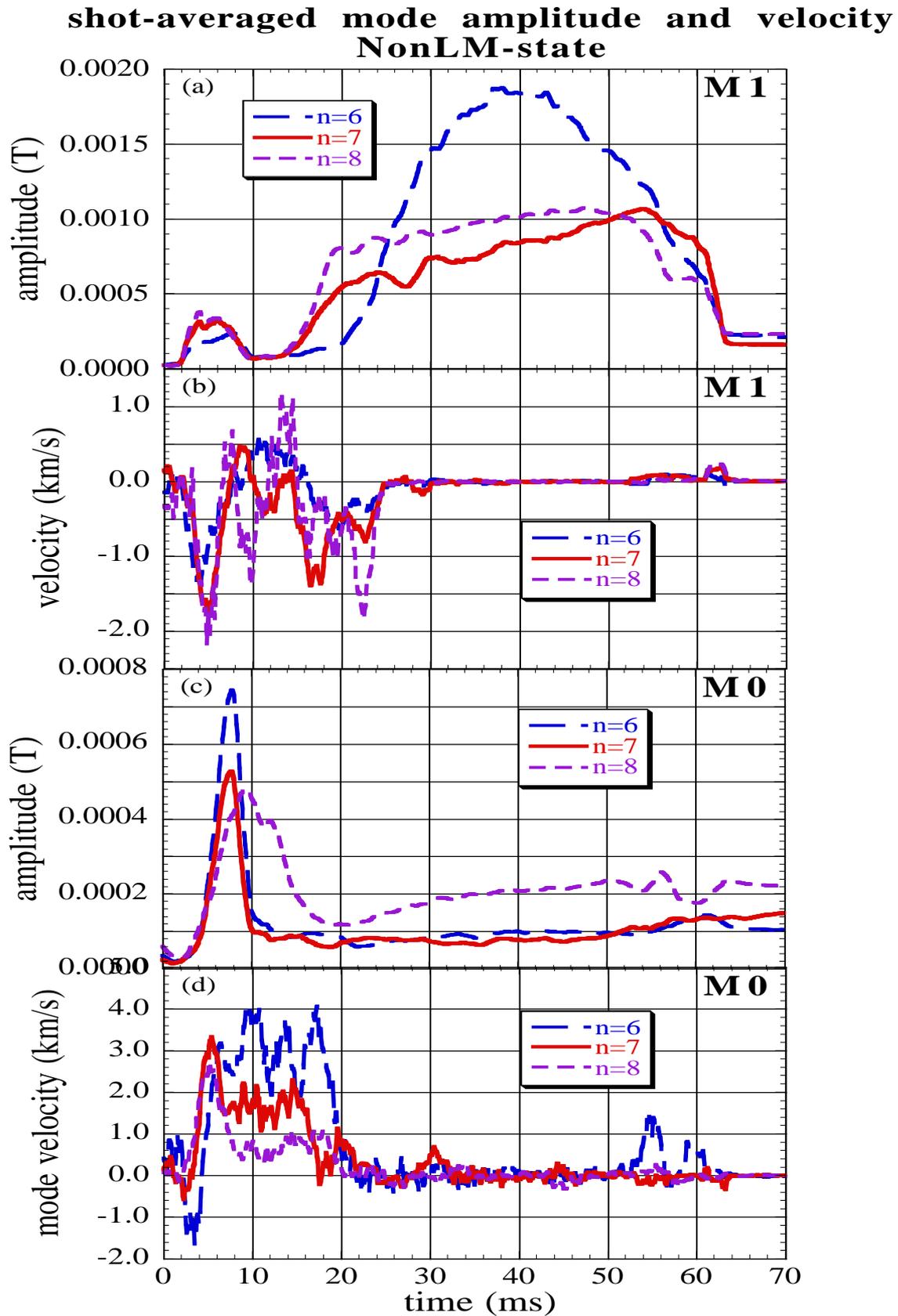


Mode Amplitude and Velocity (LM)





Mode Amplitude and Velocity (NonLM)

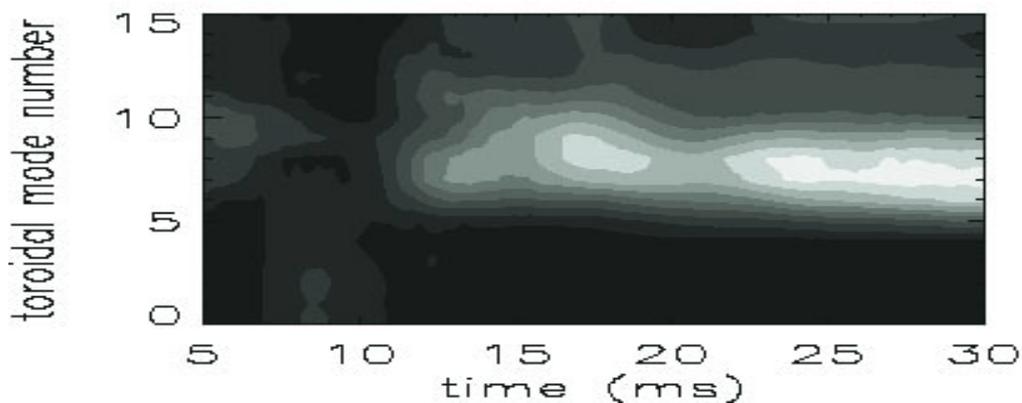




Toroidal Mode Spectrum

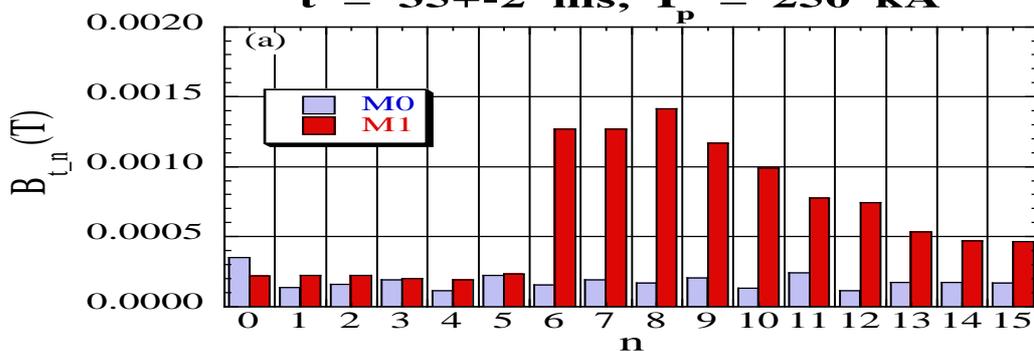
- Near-axis resonant modes ($m=1/n=8,7$) are growing in the LM case during the current- rising phase.

12888 LM-state M1 : Z=0.0299 T

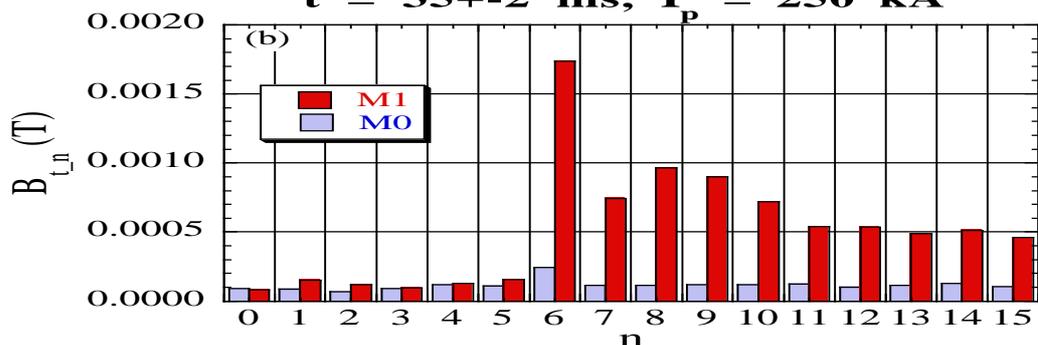


- Shot-averaged spectra at $t=35$ ms

**Ensemble average of 9 LM shots
 $t = 35 \pm 2$ ms, $I_p = 250$ kA**



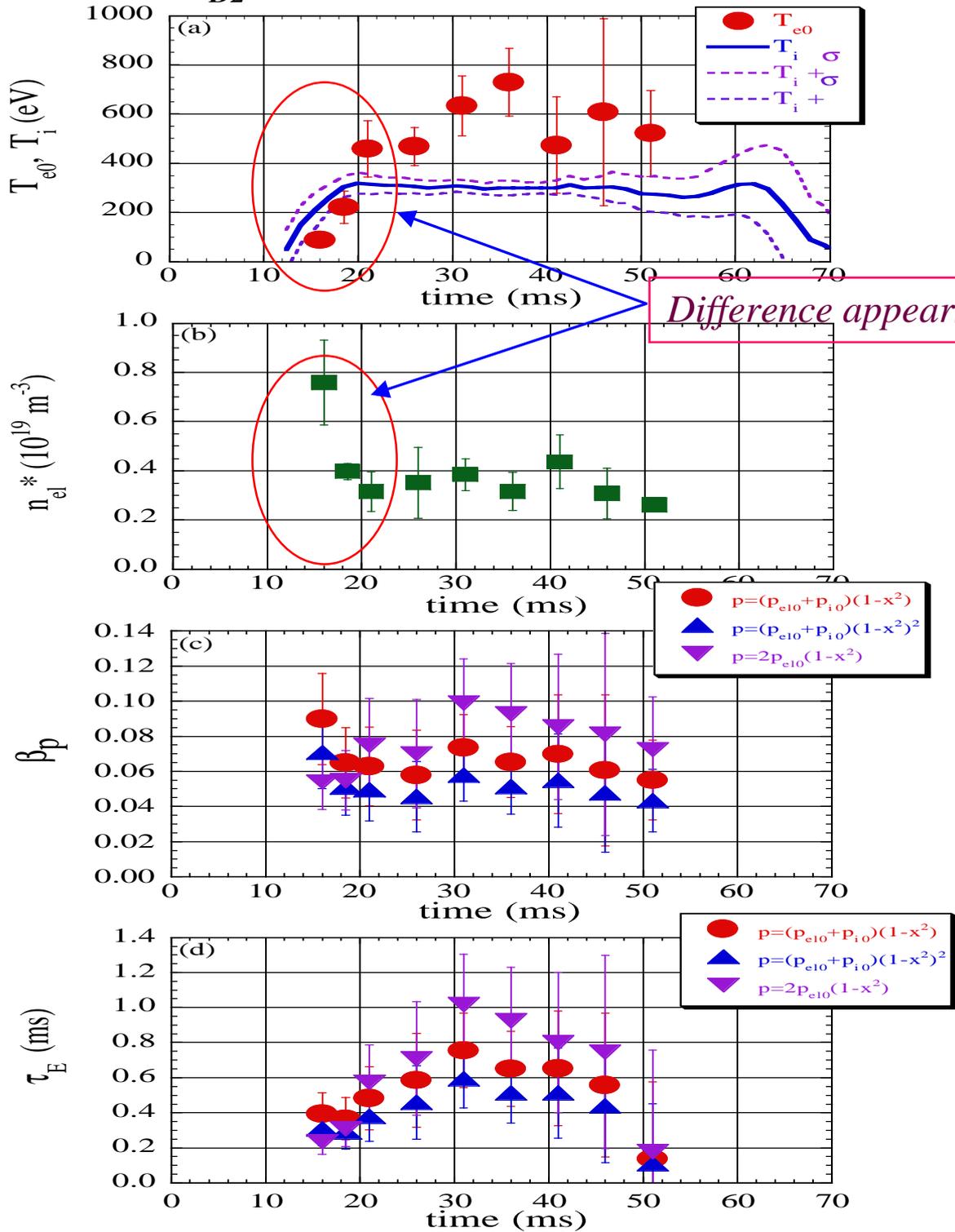
**Ensemble average of 9 NonLM shots
 $t = 35 \pm 2$ ms, $I_p = 250$ kA**





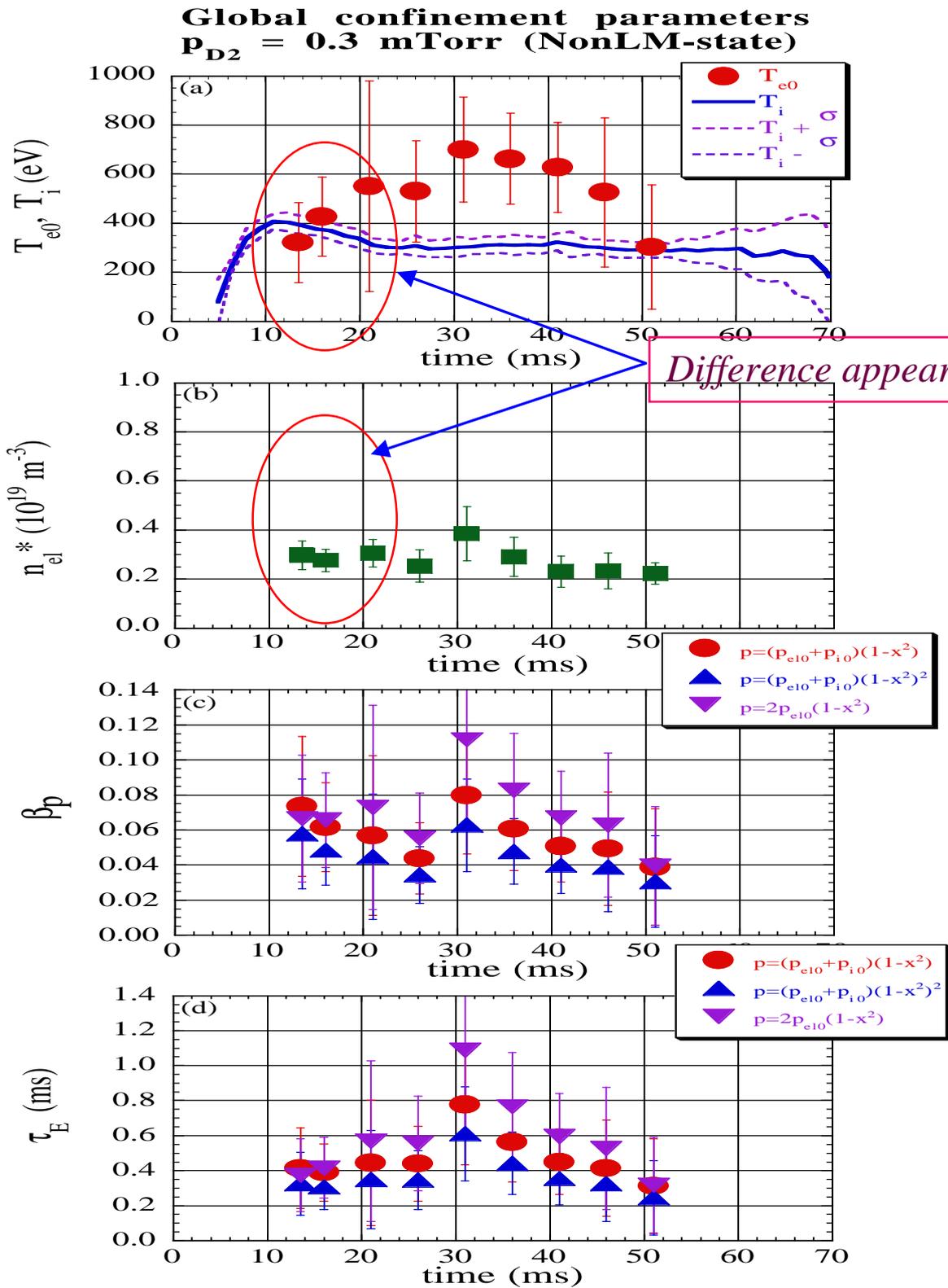
Global Confinement in LM

Global confinement parameters
 $p_{D2} = 0.7$ mTorr (LM-state)





Global Confinement in **NonLM**





Fitzpatrick's threshold

$$b_{\text{crit}}/B = f(\Theta, F) (\tau_b / \tau_f)^{1/2}$$

$$\approx T_e^{1/2} / I_p$$

when a , F , Θ , β_p and τ_E are const.

$$\tau_b = 2.5 \text{ ms} :$$

Vessel time constant

$$\tau_f = \tau_E / (\tau_A \Omega_0)^2$$

$$\tau_E = g(T_{e0}, n_{e0}, T_i, R_p, I_p) :$$

Energy confinement time

$$\tau_A = a (\mu_0 m_i n_{e0})^{0.5} / B_{pa} :$$

Alfven time

$$\Omega_0 = 6(m/n) [T_{e0} / (a^2 B_{pa})] :$$

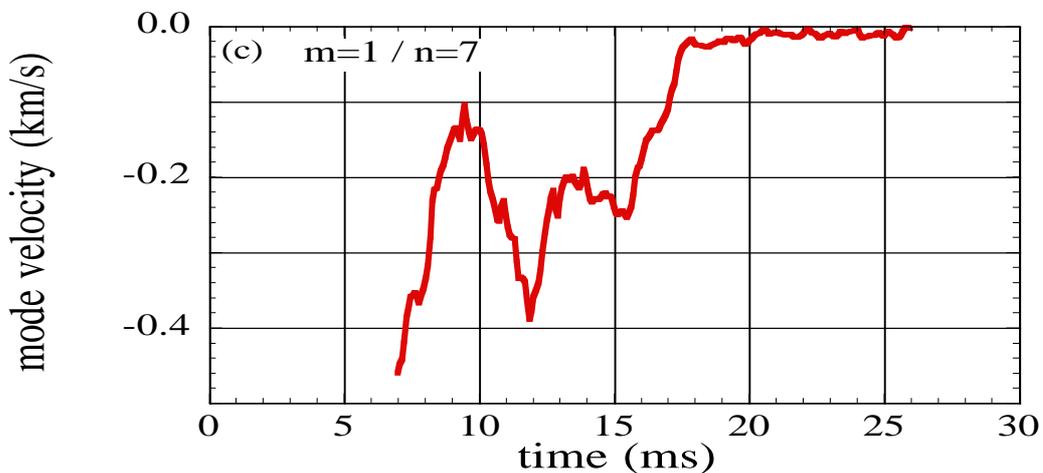
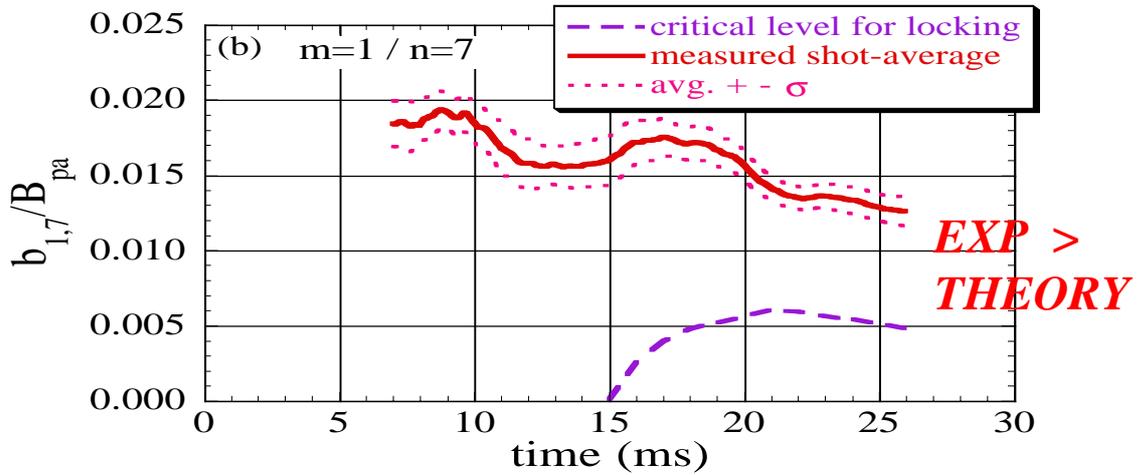
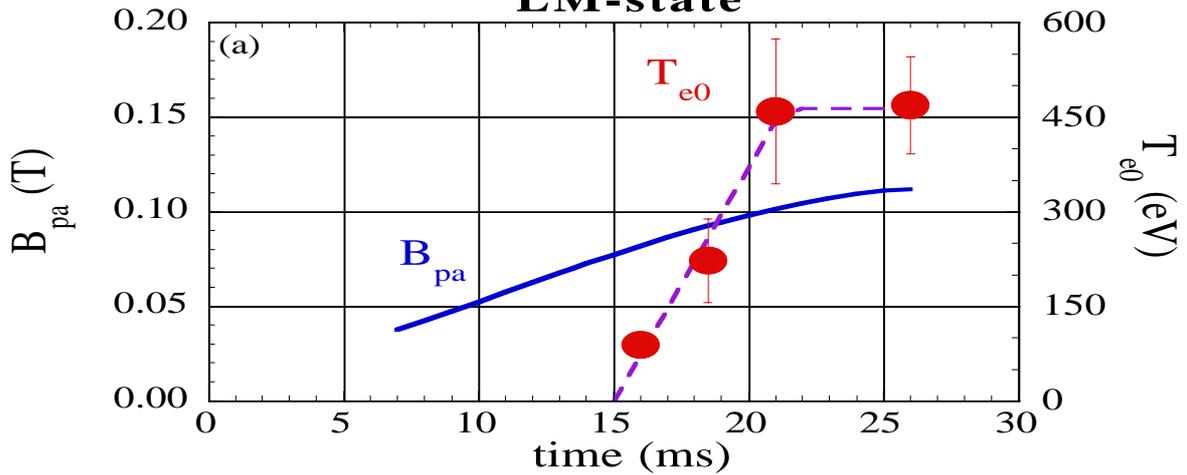
Mode rotation velocity

*Evaluation for the experimental values done by
Per Brunsell*



B_{exp}/B_{pa} and b_{theo}/B_{pa} in LM

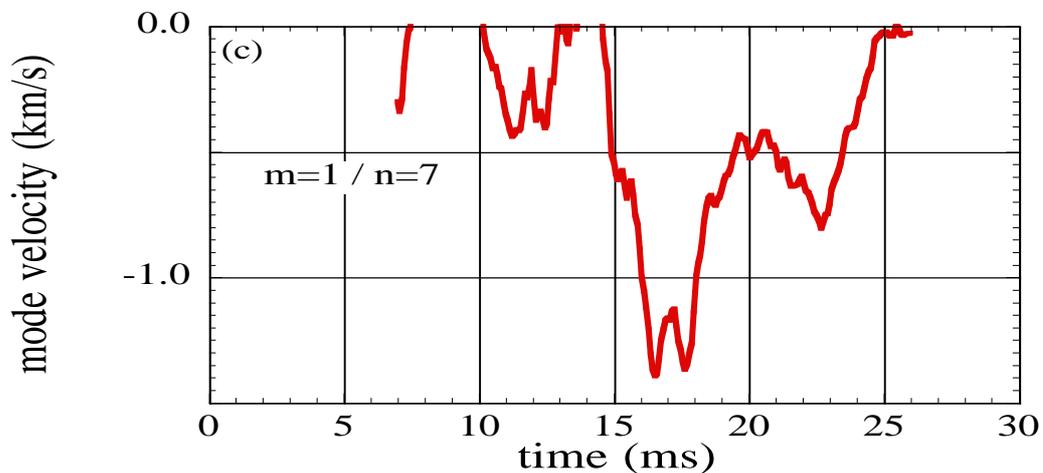
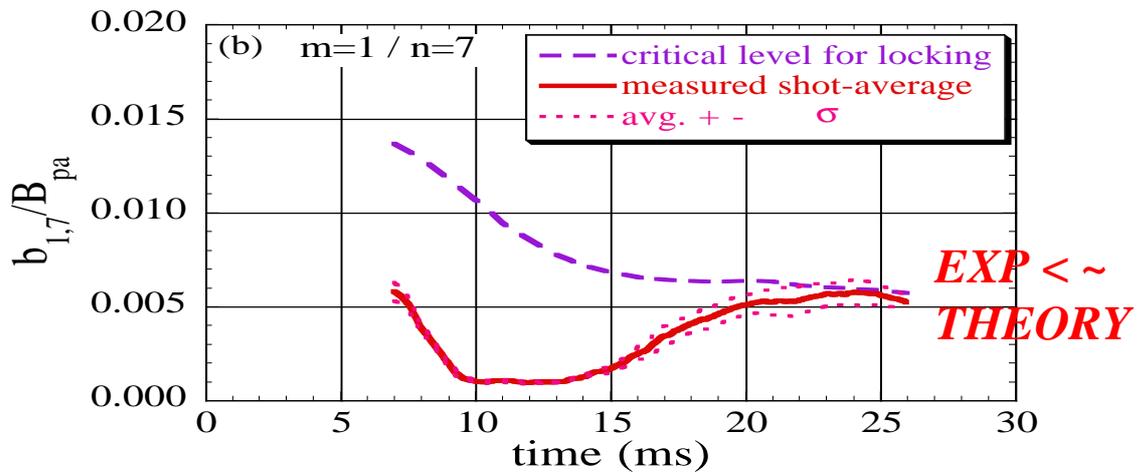
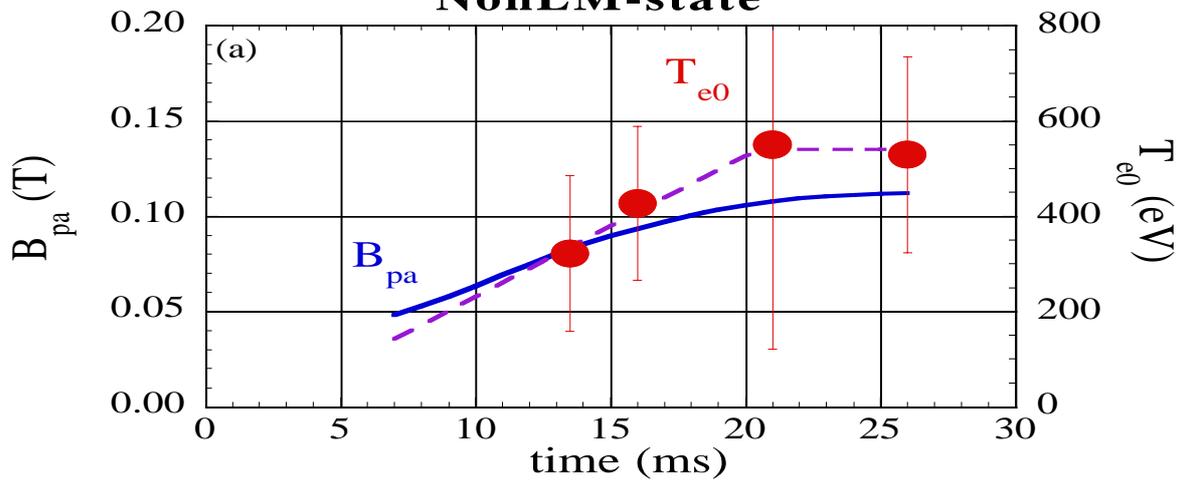
Comparison with a theoretical threshold LM-state





B_{exp}/B_{pa} and b_{theo}/B_{pa} in **NonLM**

Comparison with a theoretical threshold
NonLM-state





Summary

Phase- and wall-locked (LM) discharges are compared in detail with the nonlocked (NonLM) discharges in TPE-RX at $I_p = 250$ kA, with the B_t -array signal and the global confinement properties.

- **Rotation velocity of the modes are larger in NonLM than LM particularly during the current rising phase.**
- **Modes eventually wall-locke in NonLM but the slinky structure disappears.**
- **Global confinement properties at the center of the current flat top are comparable between LM and NonLM.**
- **T_{e0} and n_{e1}^* differ during the current rising phase, which gives rise to**
 - ***(1) a larger b (S-number related?)***
 - ***(2) a smaller theoretical threshold in the LM-state.***