



CONSORZIO RFX
Ricerca Formazione Innovazione

Recent Results About Helical States in RFX

Presented by

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at RFP Workshop Feb. 28- Mar. 1, 2000 Madison

OUTLINE

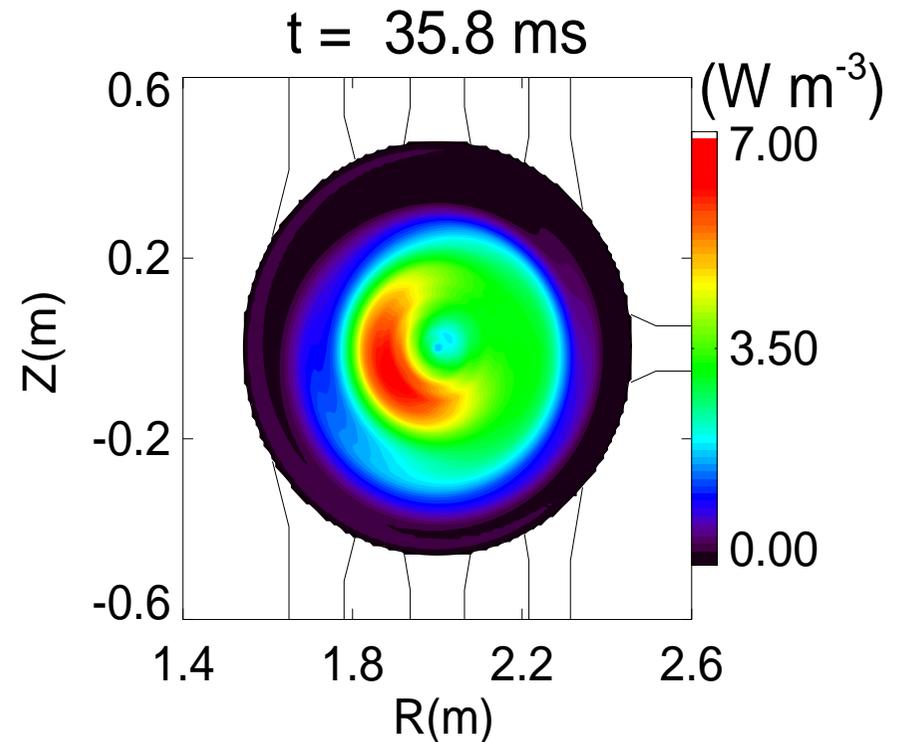
- Some experimental results on helical states in RFX
- Helical states as ohmic equilibrium states

EXPERIMENTAL EVIDENCES

- **TPE (Japan)** Brunsell et. al. (1993) : high Θ quasi single helicity states

- **T1 (Sweden)** Nordlund&Mazur (1994)

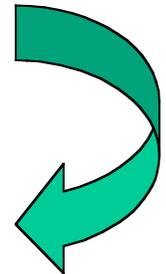
-**RFX (Italy)** : tomographic reconstruction of magnetic islands , Martin et. al. (1999)



QSH states in RFX

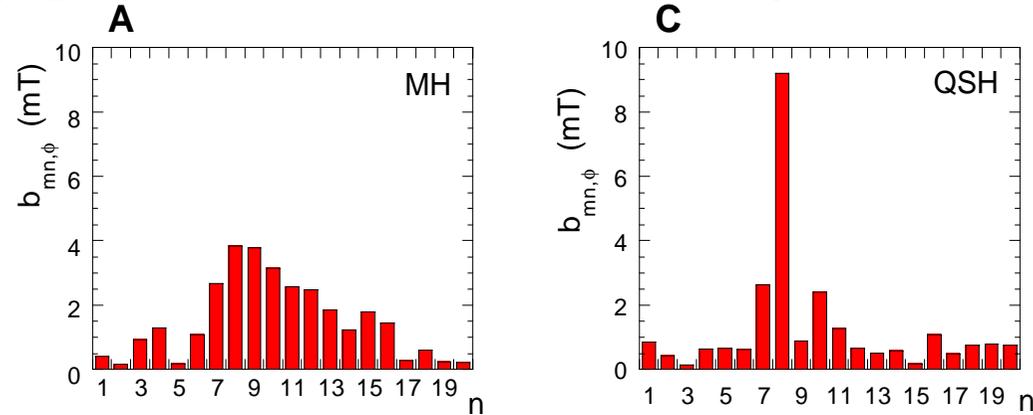
- QSH states are obtained in RFX
 - spontaneously
 - transient (similar in TPE-1RM20 - Improved High Theta Mode)
 - stationary
 - as a result of Pulsed Poloidal Current Drive (PPCD)
 - as a result of Oscillating Poloidal Current Drive (OPCD)

Techniques to control the edge
poloidal electric field

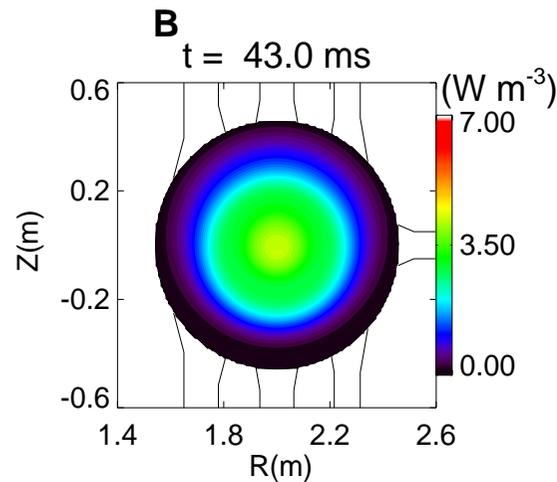


MH vs. QSH

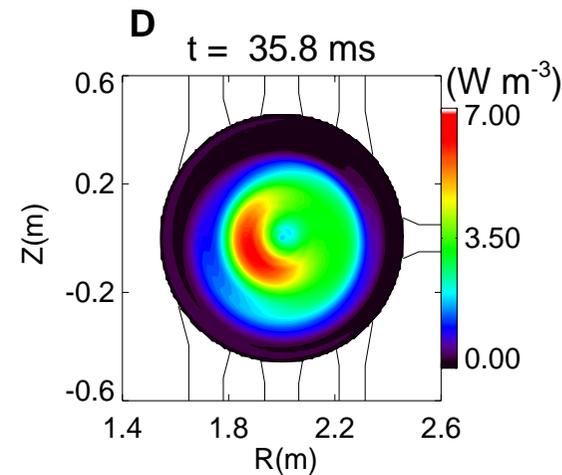
- SXR imaging reveals an helical structure in the plasma core during QSH states



MH

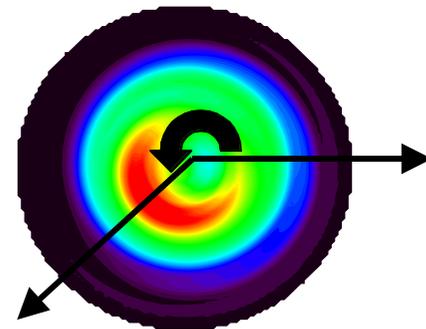
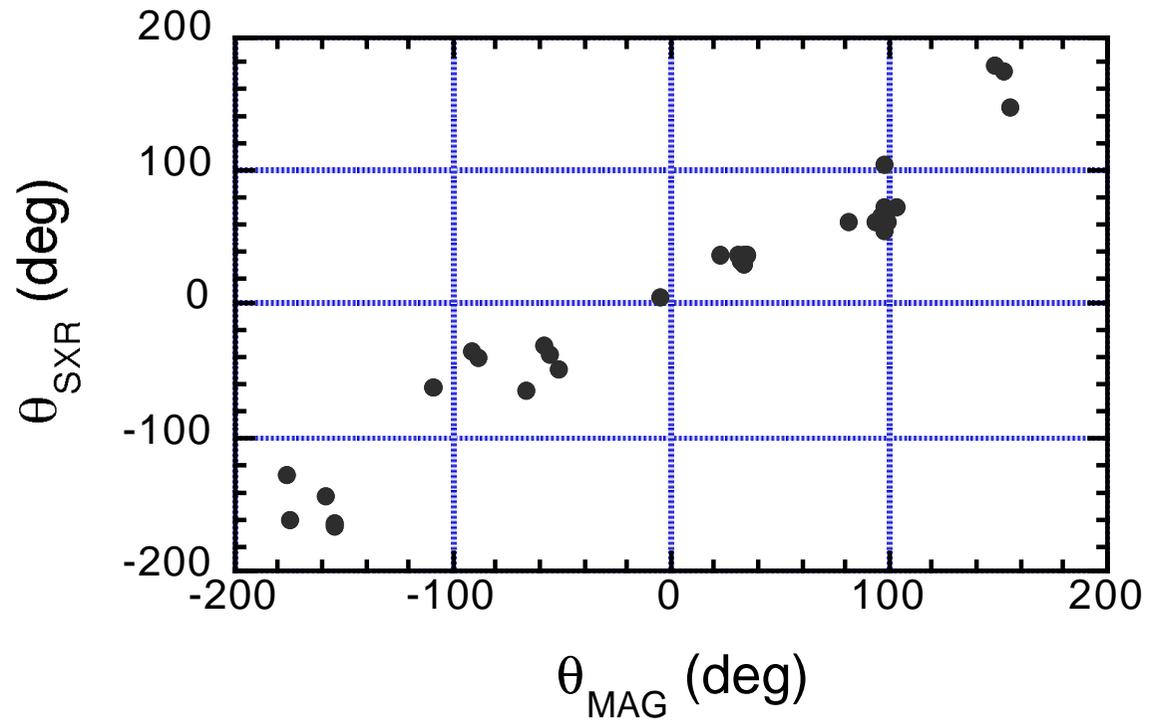


QSH

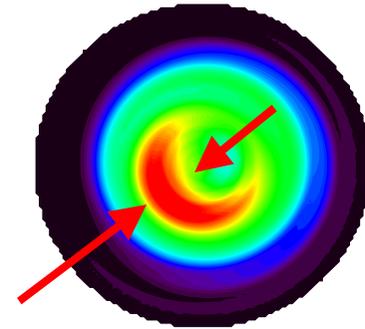
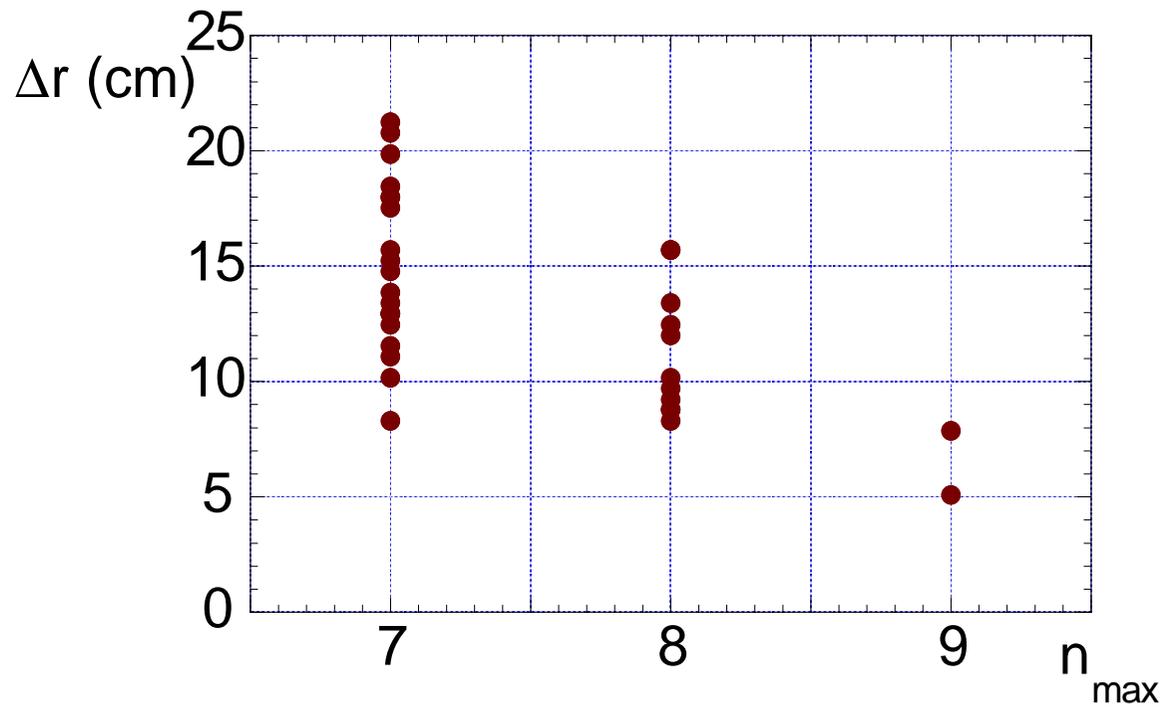


The poloidal position of the island coincides with the poloidal phase angle of the dominant mode

Poloidal position of the island

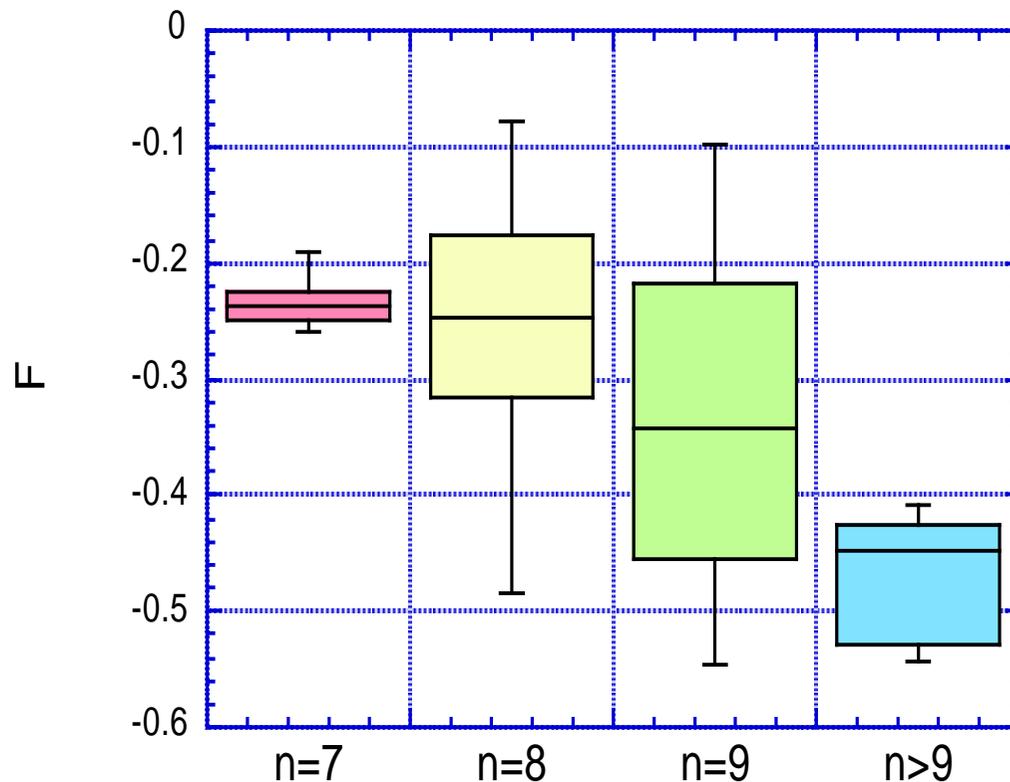


radial width of the helical structure vs. n_{\max}



For lower n the radial width of the island increases

Dominant n depends on F



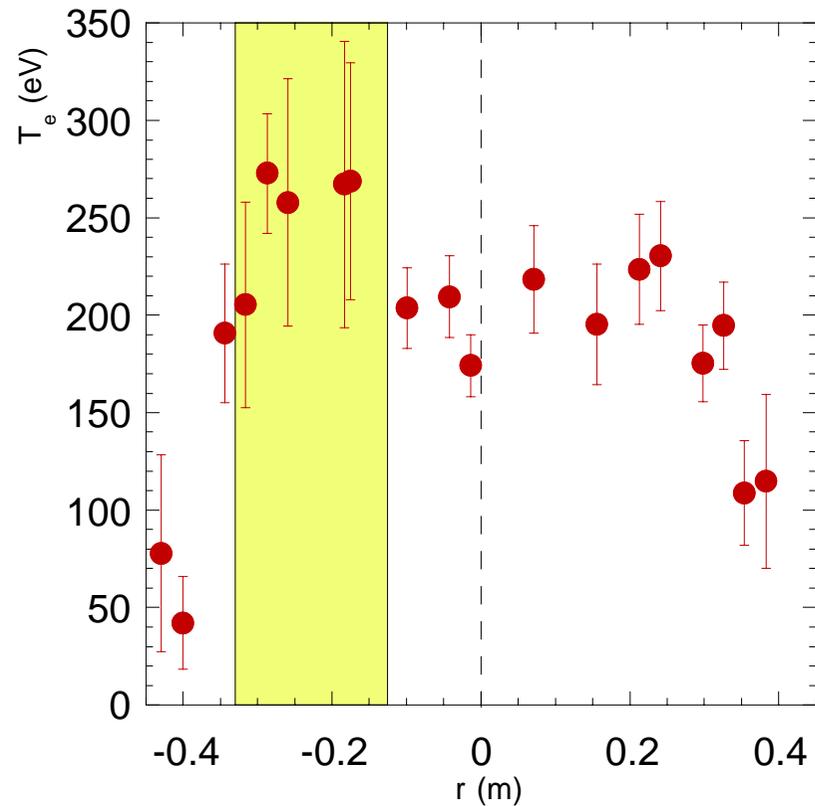
The level of the edge toroidal field reversal plays a role in determining which mode emerges in a QSH state.

A deeper reversal corresponds in fact to a lower on-axis safety factor; this can therefore eliminate the resonance of the lowest n numbers ($n=7$ with RFX aspect ratio).

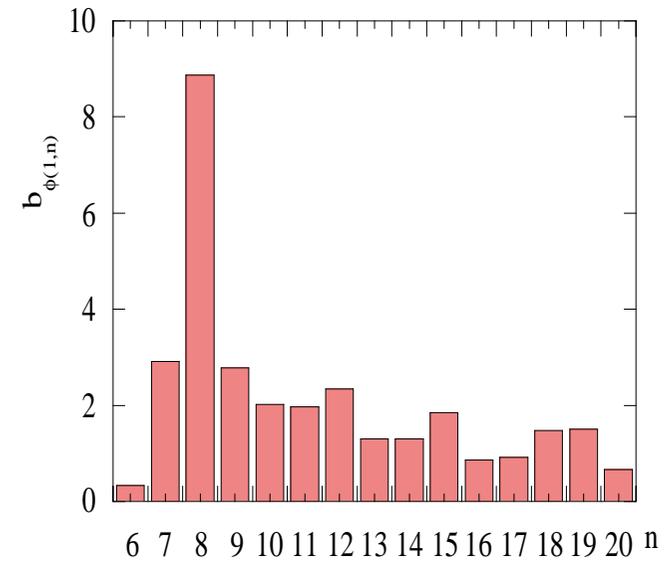
A relationship between n_{\max} and the reversal parameter F , defined as $F=B_t(a)/\langle B_t \rangle$ is indeed experimentally observed.

The core helical structure heats up

- Spontaneous QSH



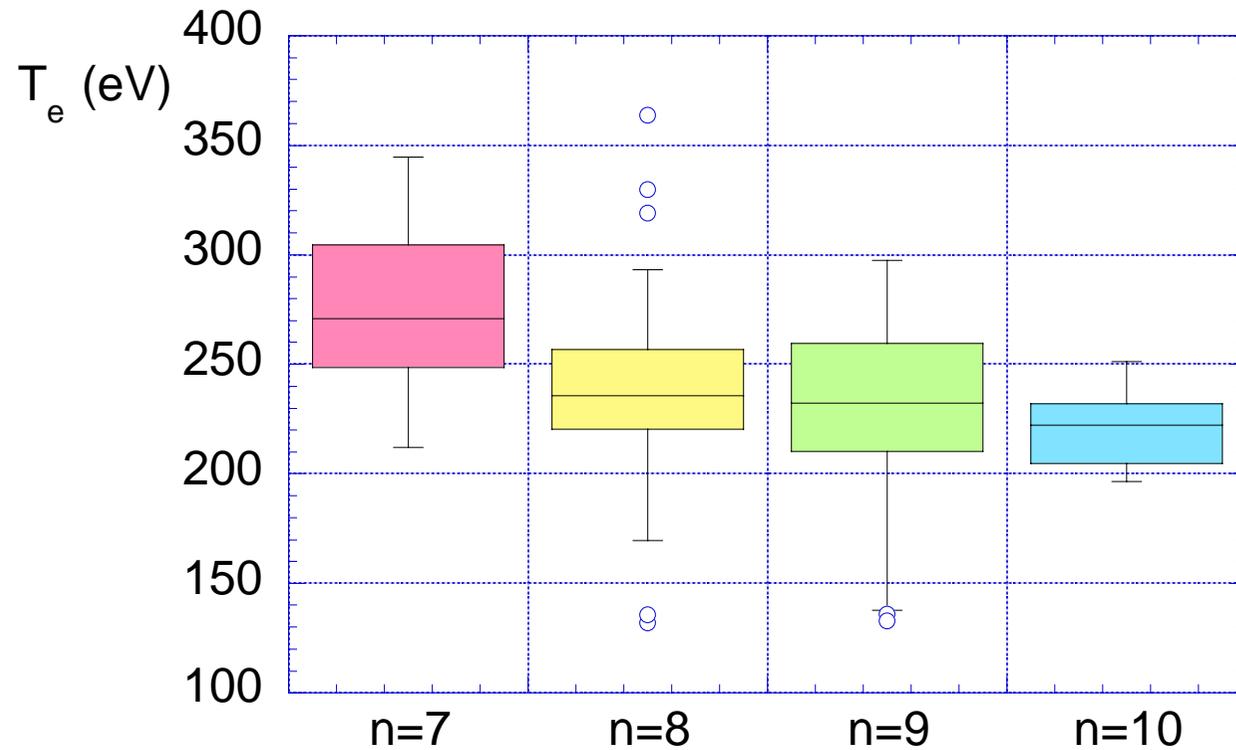
Electron temperature profile



$m=1$ toroidal spectrum

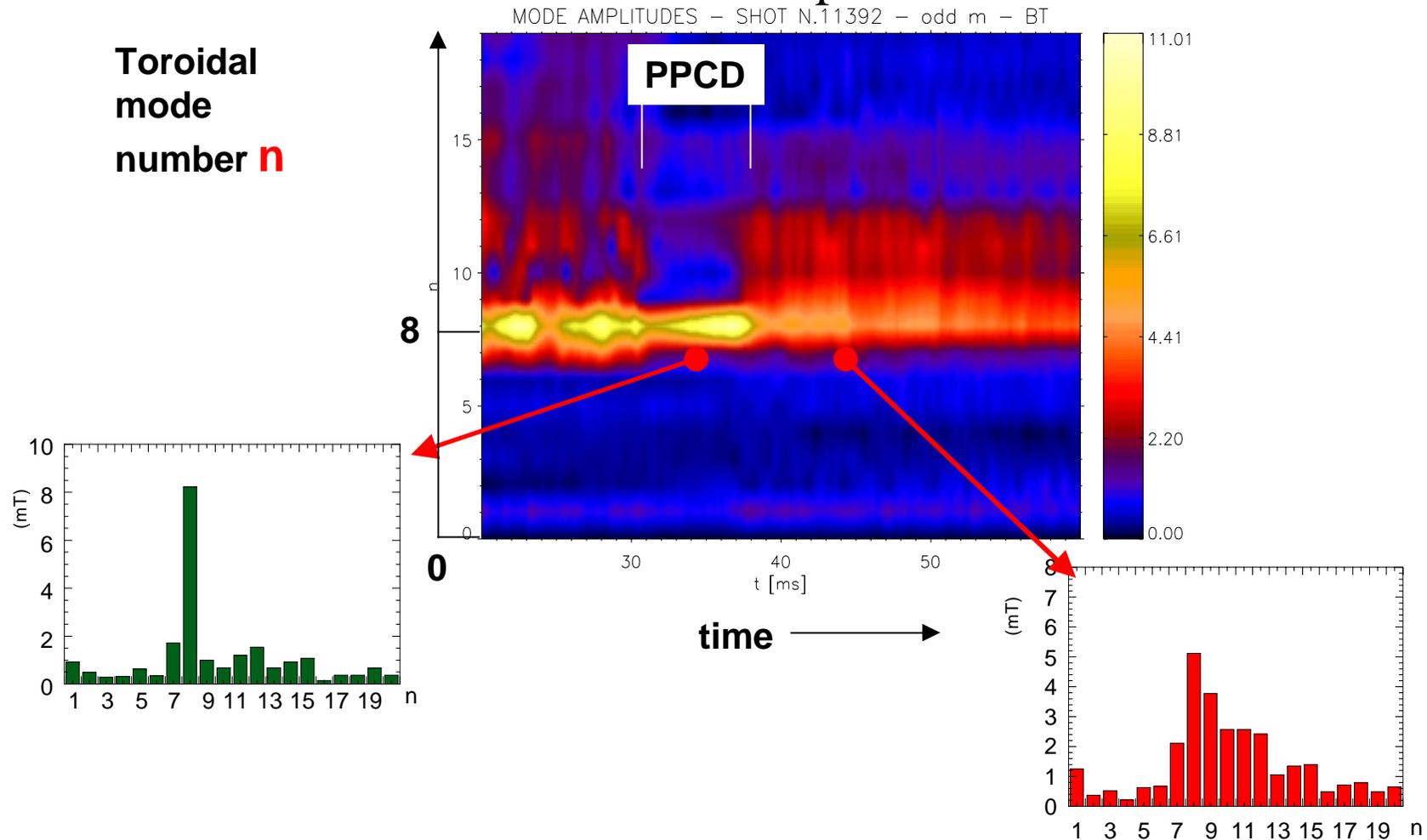
Lower n islands are hotter

- Core temperature vs. n_{dom} for QSH plasmas



Transient QSH state (during PPCD)

- Plasma makes transitions to states with a much narrower $m=1$ modes spectrum



Fluctuations and dynamo

- **A resistive, steady state, axisymmetric reversed toroidal field equilibrium is not possible**
- **RFP “dynamo” has to satisfy Cowling’s theorem:**
 - *“axisymmetric magnetic fields cannot be maintained by any symmetric fluid motion”*
- **MHD fluctuations can provide the poloidal electric field => turbulent MHD dynamo**
- **non-axisymmetric ohmic equilibrium states**

MHD Equilibrium

$$\nabla p = \mathbf{J} \times \mathbf{B}$$

Helical Grad-Shafranov Equation

$$\Delta_h \chi = \frac{1}{f} \frac{\partial}{\partial r} \left(f \frac{\partial \chi}{\partial r} \right) + \frac{1}{r f} \frac{\partial^2 \chi}{\partial u^2} = \left(c - \frac{dg}{d\chi} \right) g(\chi) - \frac{dp}{d\chi} (m^2 + k^2 r^2)$$

where χ is the helical flux function

$u = m\theta + kz$ is the helical angle, g is the helical magnetic field and p is the pressure profile

$f = r/(m^2 + (kr)^2)$, $c = -2km/(m^2 + (kr)^2)$

$\lambda(\chi) = \mathbf{J} \cdot \mathbf{B} / \mathbf{B}^2 = (dg/d\chi + g/B^2 dp/d\chi)$

OHM's Law

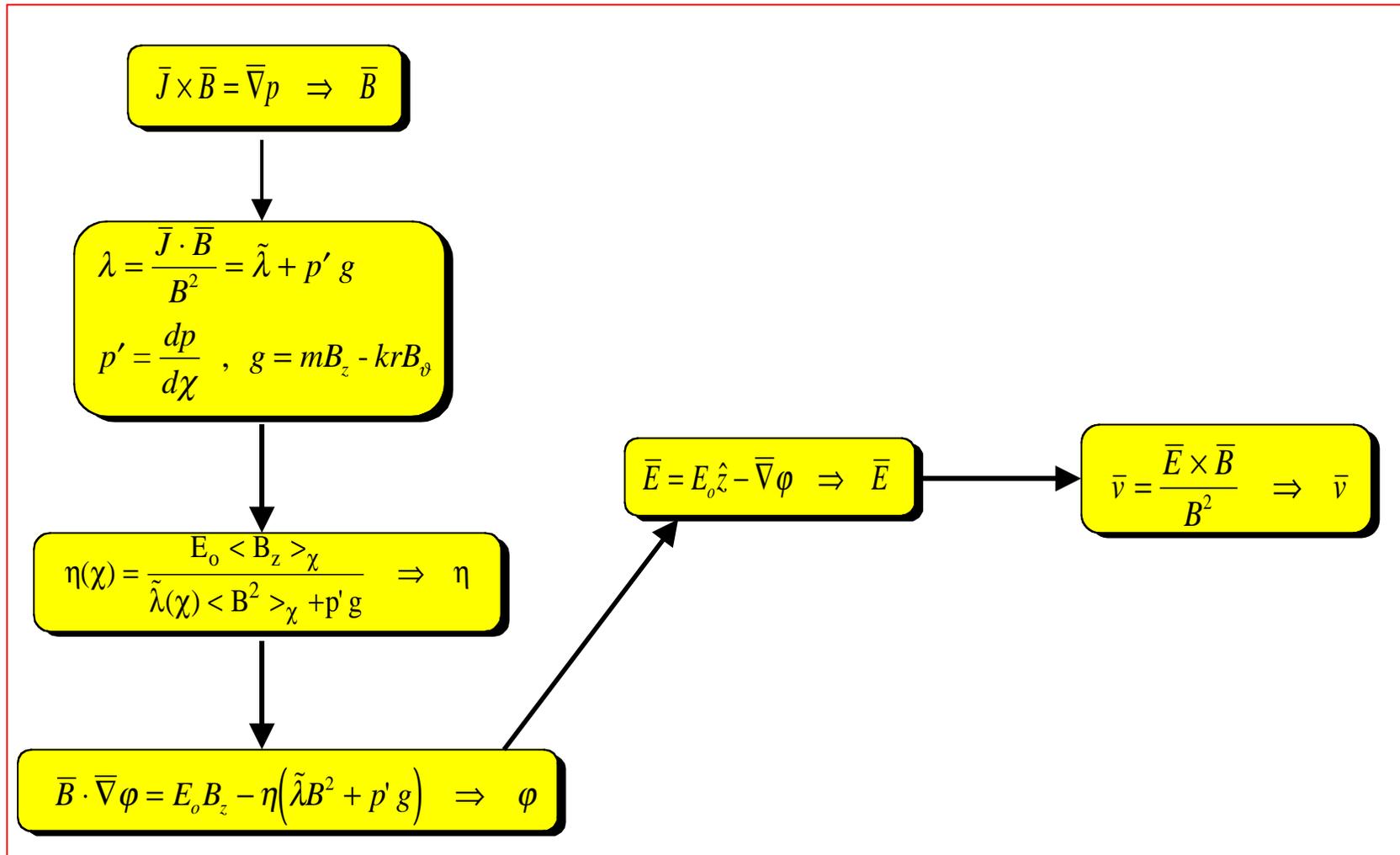
$$\vec{E} + \vec{v} \times \vec{B} = \eta \vec{j}$$

in helical symmetry
(parallel component)

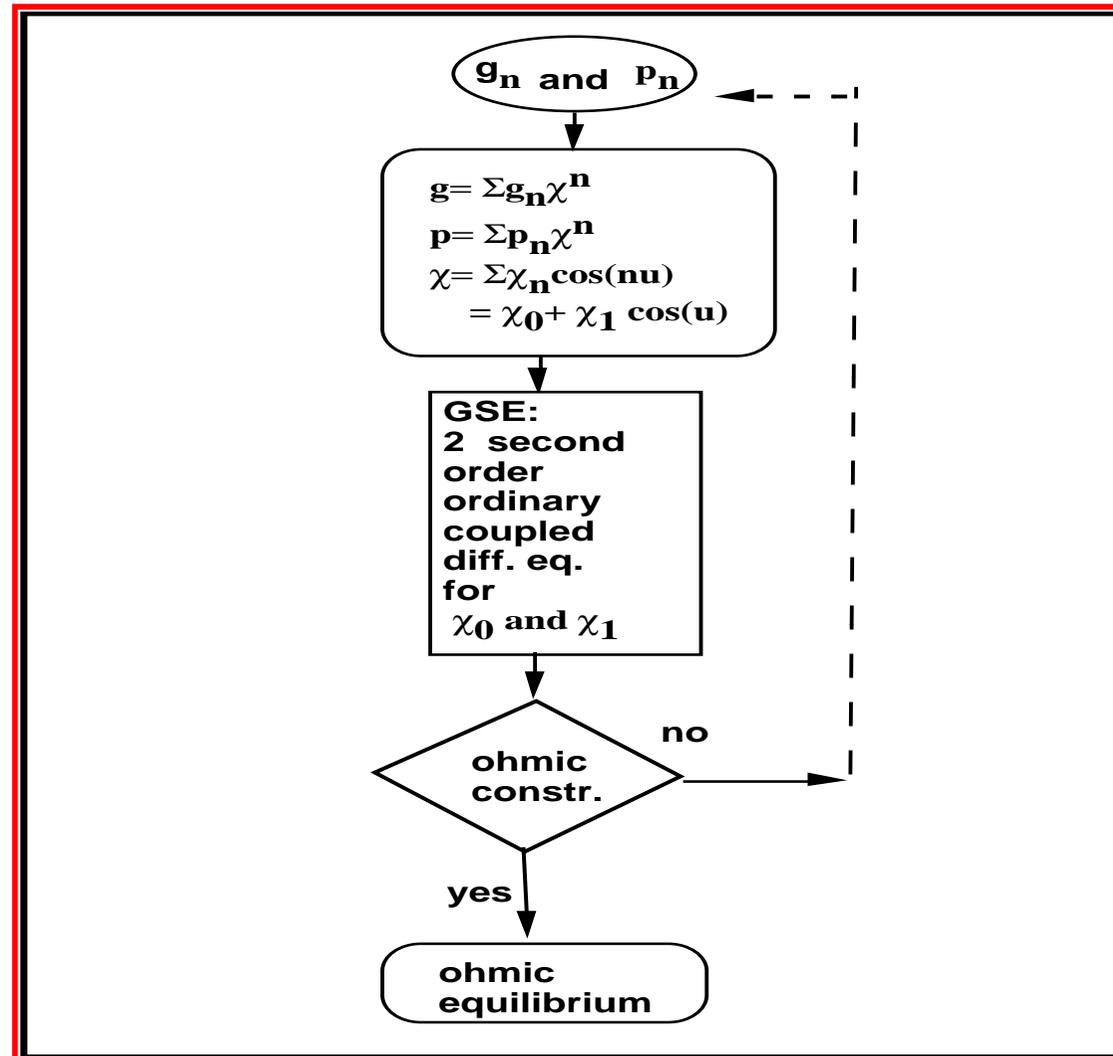
$$E_o \langle B_z \rangle_\chi = \eta(\chi) \left(\tilde{\lambda}(\chi) \langle B^2 \rangle_\chi + g \frac{dp}{d\chi} \right)$$

where $\langle \rangle$ means flux surface average and $\tilde{\lambda} = \frac{dg}{d\chi} = \vec{J} \cdot \vec{B} / B^2 - (g \mathbf{dp}/d\chi) / B^2$
Finn-Nebel-Bathke (1992)

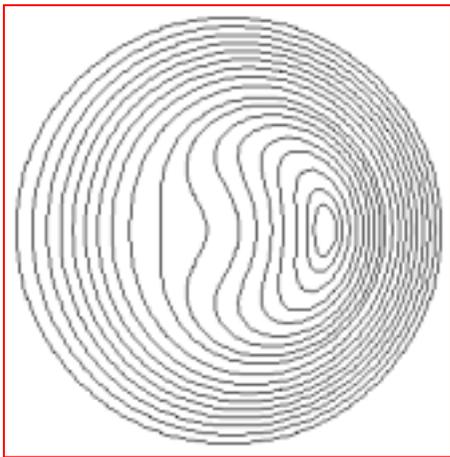
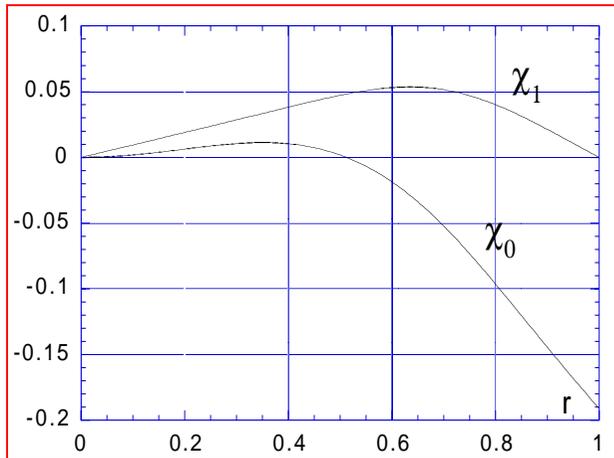
From the helical ohmic state to the calculation of the laminar velocity vector field.



Ohmic Helical equilibrium using polynomial expansion



Ohmic helical equilibrium solution



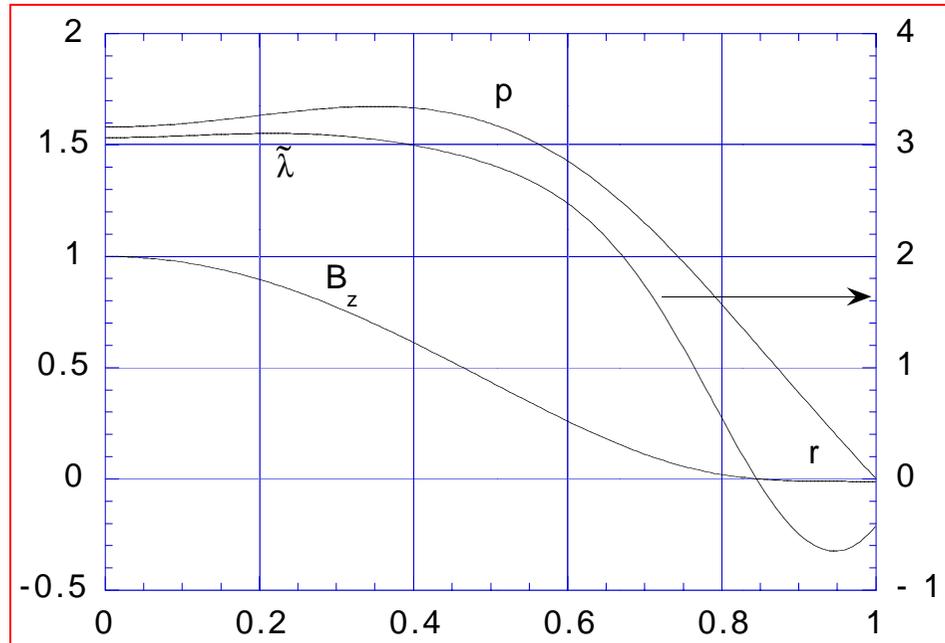
Radial profiles
and
Helical flux contours

$$\Theta=1.78$$

$$F=-0.06$$

$$\beta=0.1$$

Ohmic helical equilibrium solution (continued)

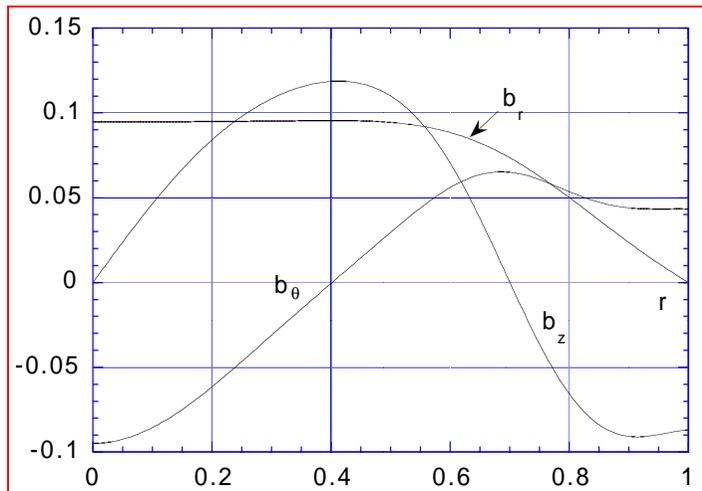
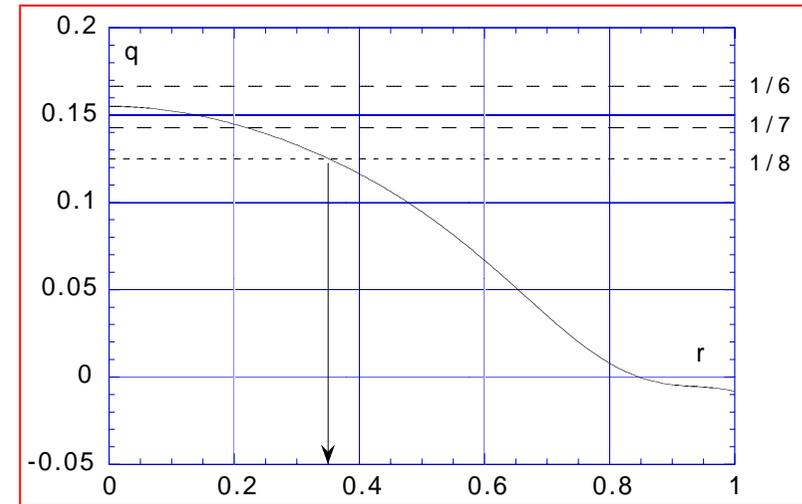
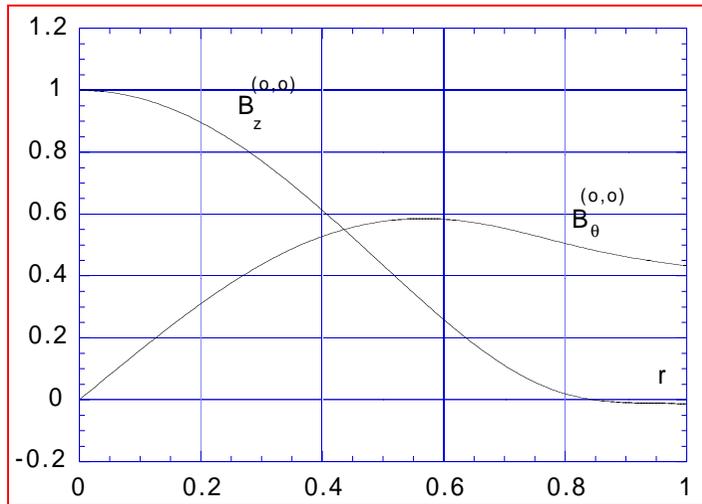


$$E_0 \langle B_z \rangle_\chi = \eta(\chi) \left(\tilde{\lambda}(\chi) \langle B^2 \rangle_\chi + g \frac{dp}{d\chi} \right)$$

with

$$\tilde{\lambda} = \frac{dg}{d\chi} = \frac{\vec{J} \cdot \vec{B}}{B^2} - (g \, d\mathbf{p}/d\chi) / B^2$$

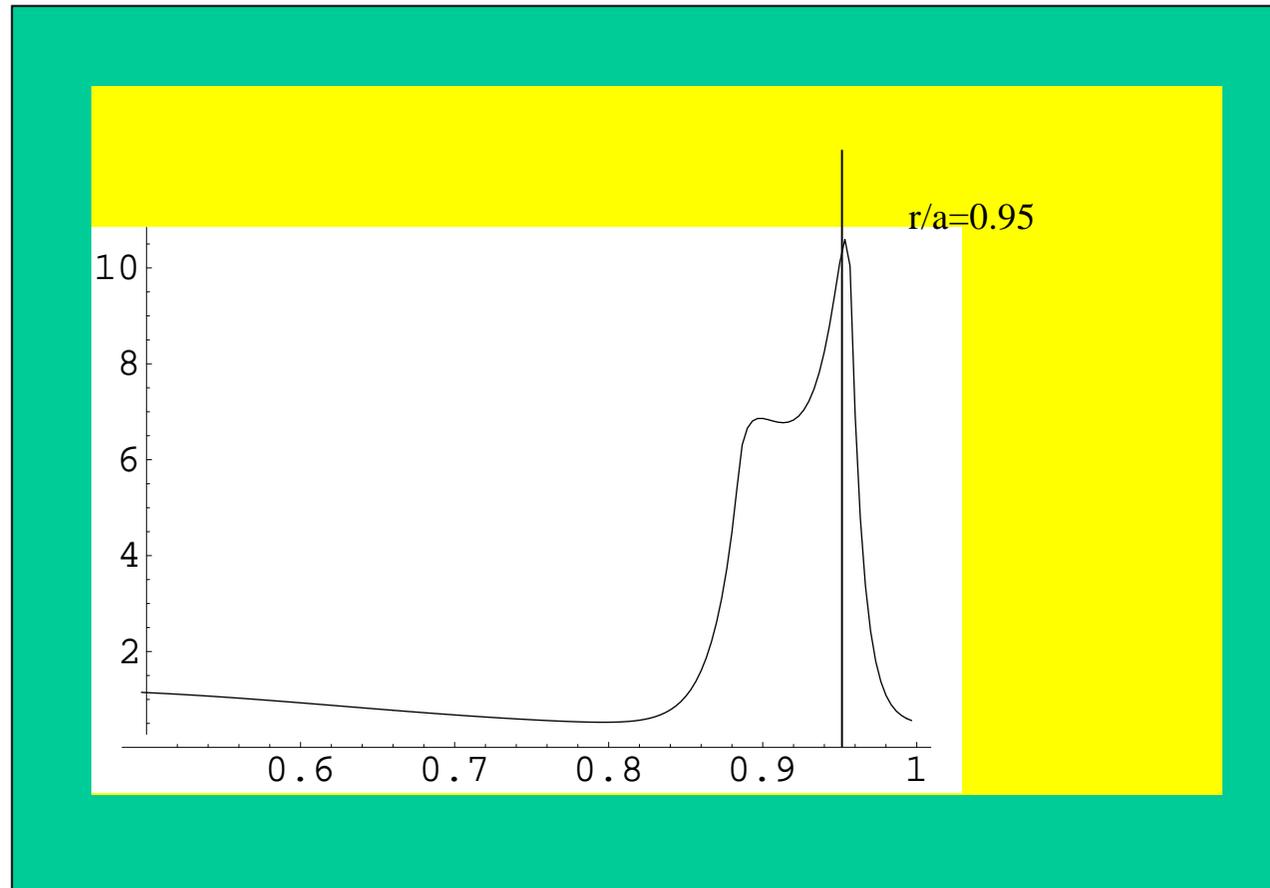
Ohmic helical equilibrium solution (continued)



Perturbation level

$$b_z(a)/B_\theta^{(0,0)}(a) \approx 20\%$$

Calculated resistivity profile



Summary

- Clear experimental evidences of helical structures in RFX
- Motivation to search ohmic helical symmetric equilibrium solutions
- Such solutions exist!
- Realistic η , b/B not too high (?)